

drift-kinetic limit,

$$|k_{\perp} \rho| \ll 1, \quad \delta E_{\parallel} \approx 0, \quad \delta B_{\parallel} \approx 0$$

$$\underline{\text{ideal MHD}} \quad \underline{\beta \ll 1}$$

$$\delta \hat{f} = -\frac{q}{m} \frac{\hat{\omega}_s}{\omega} \delta \phi F_0 + \delta \hat{H}(\varepsilon, \mu, \ell)$$

$$(2\varepsilon + v_{\perp d} + i \underline{k}_{\perp} \cdot \underline{v}_d) \delta \hat{H} = i \frac{\varepsilon}{m} (\omega \frac{\partial}{\partial \varepsilon} + \hat{\omega}_s) F_0 \frac{\omega_d}{\omega} \delta \phi$$

$\underline{v}_d$  = magnetic  $\nabla B$  and curvature drift

$$\varepsilon = v^2/2, \quad \mu = v_{\perp}^2/2B$$

$$\underline{k}_{\perp} \cdot \underline{v}_d = \omega_d,$$

$$\hat{\omega}_s F_0 = \frac{\underline{k}_{\perp} \times \underline{e}_{\parallel}}{\Omega} \cdot \nabla F_0 = \text{diamagnetic drift}$$

We then have

$$\partial_t \underline{P} \underline{u} + \underline{\nabla} \cdot \underline{P} - \underline{J} \times \underline{B} / c = 0; \quad (3)$$

where  $\underline{P} \underline{u} = \sum_{j=e,i} (m_j n_j \underline{u}_j)$ ,  $\underline{J} = \sum_j q_j n_j \underline{u}_j$  is the current density.

Linearize Eq. (3) w.r.t.  $P = P_0 + \delta P$ ,  $\underline{u} = \delta \underline{u}$ ,  $\underline{J} = \underline{J}_0 + \delta \underline{J}$ ,

$\underline{P} = P_0 \underline{I} + \delta \underline{P}$  (i.e. equilibrium pressure isotropic), we have

$$\partial_t P_0 \delta \underline{u} + \underline{\nabla} \cdot \delta \underline{P} - \frac{1}{c} (\delta \underline{J} \times \underline{B}_0 + \underline{J}_0 \times \delta \underline{B}) = 0. \quad (4)$$

Letting  $\delta \underline{u} = \partial_t \delta \underline{\xi}$  and  $\delta \underline{P} = \delta \underline{\xi} e^{-i\omega t}$ , we have

$$\boxed{\omega^2 P_0 \delta \underline{\xi} = \underline{\nabla} \cdot \delta \underline{P} - \frac{1}{c} (\delta \underline{J} \times \underline{B}_0 + \underline{J}_0 \times \delta \underline{B})} \quad (5)$$

Now  $\delta P_j = m_j \int \underline{v} \delta f_j d^3 v$ , and, from earlier notes,

$$\delta f_j = -\delta \underline{\xi} \cdot \underline{\nabla} F_{0j} + \delta H_{tj} \quad t: \text{trapped.}$$

$\delta H_{\text{circulating}} = 0.$

Hence,

$$\delta \underline{P}_j = -(\delta \underline{\xi} \cdot \underline{\nabla} P_0) \underline{I} + \sum_j m_j \langle \underline{v} \delta H_{tj} \rangle$$

$$\boxed{\equiv \delta P_{\text{conv}} \underline{I} + \delta P_{\text{comp}, t}} \quad (6)$$

Conv = convective, Comp = compressional

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Now, since  $\int d^3x = \int d\alpha \int v_1 dv_2 dv_3$ , and  $\frac{\partial \hat{H}_E}{\partial \alpha} = 0$ , thus,

it is easy to see that  $\hat{P}_{\text{comp}}$  is diagonal i.e.,

$$\hat{P}_{\text{comp},t} = \begin{bmatrix} \tilde{\delta P}_\perp & 0 \\ 0 & \tilde{\delta P}_\parallel \\ 0 & \tilde{\delta P}_\parallel \end{bmatrix} = \tilde{\delta P}_\perp \underline{\underline{I}} + (\tilde{\delta P}_\parallel - \tilde{\delta P}_\perp) \underline{\underline{e}}_1 \underline{\underline{e}}_1$$

$$\odot \frac{1}{2} \int \delta \hat{\underline{\underline{S}}}_\perp d^3x \cdot (E_2.) (5)$$

$$\Rightarrow \boxed{\omega^2 \delta \hat{\underline{\underline{I}}} = \delta \hat{W}_f + 4 \delta \hat{W}_k}, \text{ where} \quad (7)$$

$$\cdot \delta \hat{\underline{\underline{I}}} = \frac{1}{2} \int \delta \hat{\underline{\underline{S}}}_\perp^* \rho_0 d^3x$$

$$\cdot \delta \hat{W}_f = \frac{1}{2} \int d^3x \delta \hat{\underline{\underline{S}}}_\perp^* \cdot \left[ \underline{\underline{\nabla}} \delta \hat{P}_{\text{conv}} - \frac{1}{c} (\underline{\underline{J}} \times \underline{\underline{B}}_0 + \underline{\underline{J}}_0 \times \underline{\underline{\delta B}}) \right]$$

Note that  $\delta \hat{P}$ ,  $\delta \hat{\underline{\underline{J}}} = \underline{\underline{\nabla}} \times \delta \hat{\underline{\underline{B}}}$ ,  $\delta \hat{\underline{\underline{B}}} = \underline{\underline{\nabla}} \times [\delta \hat{\underline{\underline{S}}}_\perp \times \underline{\underline{B}}_0]$  are all functional of  $\delta \hat{\underline{\underline{S}}}_\perp$ . Thus,  $\delta \hat{W}_f = \delta \hat{W}_f [\delta \hat{\underline{\underline{S}}}_\perp]$ . [In fact,

one can prove that  $\delta \hat{W}_f$  is self-adjoint. one can identify it as the incompressible limit of the ideal

MHD energy principle of Bernstein et al, (1958)]

• Finally,

$$\delta \hat{W}_K = \frac{1}{2} \int d^3x \delta \hat{\underline{\underline{S}}}_2^{\wedge \#} \cdot \left[ \underline{\underline{\nabla}} \cdot \delta \hat{\underline{\underline{P}}}_2^{\wedge \#} \text{comp}, t \right]$$

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$$\underline{\underline{\nabla}} \cdot \delta \hat{\underline{\underline{P}}}_2^{\wedge \#} \text{comp}, t = \underline{\underline{\nabla}} \cdot \left[ \delta \hat{\underline{\underline{P}}}_2 \underline{\underline{I}} + (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \underline{\underline{e}}_y \underline{\underline{e}}_y \right]$$

$$= \underline{\underline{\nabla}} \delta \hat{\underline{\underline{P}}}_2 + \underline{\underline{\nabla}} \cdot (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \underline{\underline{e}}_y \underline{\underline{e}}_y$$

$$= \underline{\underline{\nabla}} \delta \hat{\underline{\underline{P}}}_2 + \underline{\underline{e}}_y \left[ \underline{\underline{\nabla}} \cdot (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \underline{\underline{e}}_y \right] + (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \underbrace{(\underline{\underline{e}}_y \cdot \underline{\underline{\nabla}}) \underline{\underline{e}}_y}_{\text{curvature}}$$

$$= \underline{\underline{\nabla}} \delta \hat{\underline{\underline{P}}}_2 + (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \left( \underline{\underline{\alpha}} + \underline{\underline{e}}_y \left[ \underline{\underline{\nabla}} \cdot (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \underline{\underline{e}}_y \right] \right)$$

$$\Rightarrow \delta \hat{W}_K^{\wedge \#} = \frac{1}{2} \int d^3x \left[ \delta \hat{\underline{\underline{S}}}_2^{\wedge \#} \cdot \underline{\underline{\nabla}} \delta \hat{\underline{\underline{P}}}_2 + (\underline{\underline{\alpha}} \cdot \delta \hat{\underline{\underline{S}}}_2^{\wedge \#}) (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) \right]$$

integrating  
by parts

$$= -\frac{1}{2} \int d^3x \left[ \delta \hat{\underline{\underline{P}}}_2 \underline{\underline{\nabla}} \cdot \delta \hat{\underline{\underline{S}}}_2^{\wedge \#} - (\delta \hat{\underline{\underline{P}}}_1 - \delta \hat{\underline{\underline{P}}}_2) (\delta \hat{\underline{\underline{S}}}_2^{\wedge \#} \cdot \underline{\underline{\alpha}}) \right]$$

as give in p. 49